

## Solutions to Problems on Assignment 9

**4.4.2(c)** The PDE we have to solve is

$$\rho_0 \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} + \alpha u,$$

where  $\rho_0$ ,  $T_0$ ,  $\alpha$  are constants, and  $\alpha < 0$ .

First find separated solutions which satisfy the boundary conditions: putting  $u(x, t) = \phi(x)h(t)$ , into the PDE and dividing by  $\rho_0\phi h$ , we get

$$\frac{\frac{d^2 h}{dt^2}}{h} = \frac{T_0}{\rho_0} \frac{\frac{d^2 \phi}{dx^2}}{\phi} + \frac{\alpha}{\rho_0}.$$

Since the left side is a function of  $t$  alone and the right side is a function of  $x$  alone, both sides must equal a constant. Let's call the constant  $-\gamma$  (so the ODE for  $h$  will have a familiar form). Then we get the ODEs

$$\frac{d^2 h}{dt^2} = -\gamma h$$

and

$$\frac{d^2 \phi}{dx^2} = \left[ \left( -\gamma - \frac{\alpha}{\rho_0} \right) \frac{\rho_0}{T_0} \right] \phi.$$

The latter equation is just of the form  $\frac{d^2 \phi}{dx^2} = (\text{constant})\phi$ , which we already know how to solve, except we usually call the constant  $-\lambda$ . So let's call the constant  $-\lambda$  here too; i.e.,

$$-\lambda = \left[ \left( -\gamma - \frac{\alpha}{\rho_0} \right) \frac{\rho_0}{T_0} \right].$$

As we know, the solutions for  $\phi$  (with boundary conditions  $\phi(0) = \phi(L) = 0$ ) are given by  $\phi(x) = \sin(\sqrt{\lambda}x)$  where  $\lambda = (n\pi/L)^2$ , for  $n = 1, 2, 3, \dots$ . Putting these values of  $\lambda$  into the preceding equation and solving for  $\gamma$ , we get

$$\gamma = \frac{(n\pi/L)^2 T_0 - \alpha}{\rho_0}.$$

Now solving the ODE for  $h(t)$  gives

$$h(t) = A \cos(\sqrt{\gamma}t) + B \sin(\sqrt{\gamma}t),$$

so the separated solutions are of the form

$$(A \cos(\sqrt{\gamma}t) + B \sin(\sqrt{\gamma}t)) \sin(\sqrt{\lambda}x),$$

where  $\lambda$  and  $\gamma$  are given by the above formulas.

Finally we have to find a superposition of the separated solutions which satisfies the correct initial conditions. Put

$$u(x, t) = \sum_{n=1}^{\infty} (A_n \cos(\sqrt{\gamma_n}t) + B_n \sin(\sqrt{\gamma_n}t)) \sin((n\pi/L)x),$$

where

$$\gamma_n = \frac{(n\pi/L)^2 T_0 - \alpha}{\rho_0}.$$

(I've attached a subscript  $n$  to the letter  $\gamma$  to emphasize that it depends on  $n$ . Notice that what I am calling  $\gamma_n$  here is called  $\lambda_n$  in the answer to this problem in the back of the book.)

Since  $u(x, 0) = 0$ , then

$$0 = \sum_{n=1}^{\infty} A_n \sin((n\pi/L)x),$$

which tells us that  $A_n = 0$  for all  $n$ . Next, differentiating  $u$  with respect to  $t$  gives

$$\frac{\partial u}{\partial t}(x, t) = \sum_{n=1}^{\infty} B_n \sqrt{\gamma_n} \cos(\sqrt{\gamma_n} t) \sin((n\pi/L)x),$$

and putting  $\frac{\partial u}{\partial t}(x, 0) = f(x)$  gives

$$f(x) = \sum_{n=1}^{\infty} B_n \sqrt{\gamma_n} \sin((n\pi/L)x),$$

which tells us that

$$B_n = \frac{2}{\sqrt{\gamma_n} L} \int_0^L f(w) \sin((n\pi/L)w) dw.$$

Therefore the solution to the problem is

$$u(x, t) = \sum_{n=1}^{\infty} B_n \sin(\sqrt{\gamma_n} t) \sin((n\pi/L)x),$$

where  $B_n$  and  $\gamma_n$  are given by the above formulas.

As discussed in class, the number  $\sqrt{\gamma_n}$  is what is usually called the *circular frequency* of this oscillation. The term *frequency* would usually be reserved for the number  $\sqrt{\gamma_n}/(2\pi)$ .

**4.4.3(b)** As in the preceding problem we put  $u(x, t) = \phi(x)h(t)$  into the PDE and divide by  $\rho_0\phi h$  to obtain

$$\frac{\frac{d^2 h}{dt^2}}{h} = \frac{T_0}{\rho_0} \frac{\frac{d^2 \phi}{dx^2}}{\phi} - \frac{\beta}{\rho_0} \frac{\frac{dh}{dt}}{h}.$$

which we can rewrite as

$$\frac{\rho_0}{T_0} \left[ \frac{\frac{d^2 h}{dt^2}}{h} + \frac{\beta}{\rho_0} \frac{\frac{dh}{dt}}{h} \right] = \frac{\frac{d^2 \phi}{dx^2}}{\phi}.$$

Setting both sides equal to the constant  $-\lambda$ , we get for  $\phi$  the usual ODE  $\frac{d^2 \phi}{dx^2} = -\lambda\phi$  and boundary conditions  $\phi(0) = \phi(L) = 0$ . Thus  $\lambda = -(n\pi/L)^2$  and  $\phi(x) = \sin((n\pi/L)x)$ . For  $h(t)$  we get the ODE

$$\frac{d^2 h}{dt^2} + \frac{\beta}{\rho_0} \frac{dh}{dt} = - \left( \frac{\lambda T_0}{\rho_0} \right) h.$$

The characteristic equation for this ODE is

$$r^2 + \left( \frac{\beta}{\rho_0} \right) r + \left( \frac{\lambda T_0}{\rho_0} \right) = 0,$$

and using the quadratic equation to solve it gives

$$r = \frac{\left( \frac{-\beta}{\rho_0} \right) \pm \sqrt{\left( \frac{\beta}{\rho_0} \right)^2 - 4 \left( \frac{\lambda T_0}{\rho_0} \right)}}{2} = \frac{-\beta}{2\rho_0} \pm \sqrt{\frac{\beta^2}{4\rho_0^2} - \left( \frac{\lambda T_0}{\rho_0} \right)} = \frac{-\beta}{2\rho_0} \pm \sqrt{\frac{\beta^2}{4\rho_0^2} - \frac{T_0}{\rho_0} \left( \frac{n\pi}{L} \right)^2}.$$

Since the problem says we can assume that  $\beta^2 < 4\phi^2\rho_0 T_0/L^2$ , then in the last equation the quantity under the radical is negative, and so can be written as  $-w_n^2$ , where  $w_n$  is given by

$$(1) \quad w_n = \sqrt{\frac{T_0}{\rho_0} \left( \frac{n\pi}{L} \right)^2 - \frac{\beta^2}{4\rho_0^2}}.$$

We then have

$$r = -(\beta/2\rho_0) \pm iw_n.$$

The general solution of the ODE for  $h$  is now given by

$$\begin{aligned} h(t) &= Pe^{-(\beta/2\rho_0)+iw_nt} + Qe^{-(\beta/2\rho_0)-iw_nt} \\ &= e^{-\beta t/2\rho_0} (Pe^{iw_nt} + Qe^{-iw_nt}) \\ &= e^{-\beta t/2\rho_0} (A \cos(w_nt) + B \sin(w_nt)), \end{aligned}$$

where  $A$  and  $B$  are arbitrary constants.

So, separated solutions of the PDE and boundary conditions are of the form

$$e^{-\beta t/2\rho_0} [A \cos(w_nt) + B \sin(w_nt)] \sin(n\pi x/L).$$

Next, we look for a superposition

$$(2) \quad u(x, t) = \sum_{n=1}^{\infty} e^{-\beta t/2\rho_0} [A_n \cos(w_nt) + B_n \sin(w_nt)] \sin(n\pi x/L)$$

which satisfies the given initial conditions. Putting  $t = 0$  and using  $u(x, 0) = f(x)$ , we find that

$$f(x) = \sum_{n=1}^{\infty} A_n \sin(n\pi x/L),$$

so

$$(3) \quad A_n = \frac{2}{L} \int_0^L f(w) \sin(n\pi w/L) dw.$$

Also, taking the derivative of  $u(x, t)$  with respect to  $t$  gives

$$\begin{aligned} \frac{\partial u}{\partial t} &= \\ \sum_{n=1}^{\infty} \left\{ -\frac{\beta}{2\rho_0} e^{-\beta t/2\rho_0} [A_n \cos(w_nt) + B_n \sin(w_nt)] + e^{-\beta t/2\rho_0} [-A_n w_n \sin(w_nt) + B_n w_n \cos(w_nt)] \right\} \sin(n\pi x/L), \end{aligned}$$

and putting  $t = 0$  gives

$$g(x) = \sum_{n=1}^{\infty} \left[ -\left(\frac{\beta}{2\rho_0}\right) A_n + B_n w_n \right] \sin(n\pi x/L).$$

Therefore

$$\left[ -\left(\frac{\beta}{2\rho_0}\right) A_n + B_n w_n \right] = \frac{2}{L} \int_0^L g(w) \sin(n\pi w/L) dw,$$

and solving for  $B_n$  gives

$$(4) \quad B_n = \frac{2}{L w_n} \int_0^L g(w) \sin(n\pi w/L) dw + \frac{\beta A_n}{2\rho_0 w_n}.$$

The solution is therefore given by the series in equation (2), where  $w_n$  is given by equation (1),  $A_n$  is given by equation (3), and  $B_n$  is given by equation (4).